# Elastic Surface Model For Beta-Barrels: Geometric, Computational, And Statistical Analysis 

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#### Abstract

Over the past 2 decades, many different geometric models were created for beta barrels, including, but not limited to: cylinders, 1-sheeted hyperboloids, twisted hyperboloids, catenoids, and so forth. We are proponents of an elastic surface model for beta-barrels, which includes the minimal surface model as a particular case, but is a lot more comprehensive. Beta barrel models are obtained as numerical solutions of a boundary value problem, using the COMSOL Multiphysics Modeling Software. We have compared them against the best fitting statistical models, with positive results. The geometry of each individual beta barrel, as a rotational elastic surface, is determined by the ratio between the exterior diameter and the height. Through our COMSOL computational modeling, we created a rather large variety of generalized Willmore surfaces that may represent models for beta barrels. The catenoid is just a particular solution among all these shapes.


## KEYWORDS

beta barrel, generalized Willmore energy, mean curvature, minimal surface, protein structure

## 1 | INTRODUCTION

In biochemistry, biophysics and mathematical biology, secondary structures represent the main types of local surfaces (shapes) corresponding to biopolymers (eg, proteins and nucleic acids). On a finer level, the atomic positions in 3D space are said to form the tertiary structure.

The most common secondary structures are the alpha helices. The second most-common are the beta sheets and beta barrels. A beta barrel is a collection of beta-sheets that twist and coil in a shape that can be described as a smooth surface of revolution which resembles a barrel. In this structure, the first strand is hydrogen bonded to the last. Beta-strands in beta-barrels are typically arranged in an antiparallel fashion. Barrel structures are commonly found in proteins that span cell membranes and in proteins that bind hydrophobic ligands in the barrel center.

A beta barrel represents an ideal smooth surface described by the beta strands. The most common type, the so called up-and-down barrel, is considered to be a surface of revolution that is topologically equivalent to a cylinder. Up-and-down barrels consist of a series of beta strands, each of which is hydrogen-bonded to the strands immediately before and after it in the primary sequence. Beta-strands in beta-barrels are typically arranged in an antiparallel fashion, but some proteins, such as the green fluorescent protein (GFP), are characterized
by beta barrels formed with both parallel and antiparallel beta strands. Beta barrel structures (named for their resemblance to the barrels that hold liquids) are commonly found in porins and other proteins that span cell membranes, and in proteins that bind hydrophobic ligands in the barrel center, as in lipocalins.

In many cases, the strands contain alternating polar and hydrophobic amino acids, so that the hydrophobic residues are oriented toward the interior of the barrel to form a hydrophobic core and the polar residues are oriented toward the outside of the barrel on the solventexposed surface. Porins and other membrane proteins containing beta barrels reverse this pattern, with hydrophobic residues oriented toward the exterior where they contact the surrounding lipids, and hydrophilic residues oriented toward the interior pore.

The polyhedral, discrete skeleton of beta-barrels can be classified in terms of 2 integer parameters: the number of strands in the betasheet, $n$, and the "shear number", $S$, a measure of the stagger of the strands in the beta-sheet. These 2 parameters ( $n$ and $S$ ) are related to the inclination angle of the beta strands relative to the axis of the barrel. To us, the number of strands and the sheer number are irrelevant. We consider the best fitting smooth surface for this polyhedral molecular skeleton, to be what we refer to as beta barrel.

Several models were proposed for beta sheets and beta barrels. Among them, we recall, in chronological order: the twisted 1-sheeted

TABLE 1 Mean curvature values of beta barrels for several types of proteins (from reference ${ }^{4}$ )

| Protein | Average <br> of mean <br> curvature | SD of mean <br> curvature |
| :--- | :--- | :--- |
| Triose phosphate isomerase | 0.040 | 0.021 |
| Taka-amylase | 0.035 | 0.007 |
| Glycolate-oxidase | 0.035 | 0.007 |
| Trimethanolamine <br> dehydrogenase | 0.037 | 0.013 |
| Cytochrome b2 | 0.033 | 0.005 |
| Aldolase | 0.035 | 0.112 |

hyperboloid (see, ${ }^{1}$ 1984), followed by the usual 1-sheeted hyperboloid (see, ${ }^{2}$ 1988), and much later, the catenoid as a best-fit (see, ${ }^{3}$ 2005). Over time, all the above-mentioned surfaces have been tried as "best models" for beta barrels. We became aware of the fact that none of these models is satisfactory. A few authors presented arguments based on experimental data that the beta sheet structures in proteins are the result of the tendency to minimize surface areas, and should therefore be very close to minimal surfaces. For a large diversity of aminoacids, the mean curvature was experimentally measured and it turned out to be close to a specific constant, which is small in absolute value, but not negligible. A decade ago, Koh and Kim ${ }^{3}$ have proposed the model that all beta-sheet structures "are almost minimal surfaces". By this, they meant that their mean curvatures are very small, but the term almost-minimal has no mathematical foundation.

We would like to make it clear, however, that while the mean curvature could be small in absolute value for some beta barrels, it may be far away from zero for others, and it may also vary significantly from a point to another of the same barrel. Koh and Kim stated: "The fact that the commonly used models for some beta-sheet surfaces (ie, the hyperboloid and strophoid) have very small mean curvatures (under 0.05) supports our model". For example, for the following enzymes: glycolate-oxidase, taka-amylase, and aldolase, the mean curvature H , measured experimentally for beta-sheets, is approximately $H=0.039$ (for each of them). In reference, ${ }^{4}$ the authors presented the mean curvature values for several aminoacids, and are depicted in Table 1. On the other hand, the mean curvature values may change significantly from a type of protein to another.

Moving in a new direction, that can give us a better view, we are proponents of an "elastic surface model for beta-barrels", which includes the minimal surface model as a particular case, but is a lot more comprehensive. Such a model has already been proposed in part by ${ }^{5}$ for beta sheets with antiparallel strands. We study beta barrels as elastic surfaces, and obtain them as numerical solutions of a boundary value prob-lem-while comparing them against the best fitting statistical models. It is important to remark that catenoids represent the only minimal surfaces of revolution, and in particular, they are Willmore-type surfaces (elastic surfaces) [14]. As such, they are included in our model. Through our COMSOL computational modeling, we created a rather large variety of Willmore-type surfaces that may represent models for beta barrels.

As a relevant real-world application of high relevance and actuality, observe the model of the beta barrel of GFP in Figure 1 and remark its unduloidal shape (Delaunay-type unduloid). Martin Chalfie, Osamu Shimomura, and Roger Y. Tsien were awarded the 2008 Nobel Prize in Chemistry for their discovery and development of the GFP model. On the other hand, Helfrich [10] has introduced the curvature energy per unit area, corresponding to bio-membranes (lipid bilayers) as:

$$
\begin{equation*}
E_{\mathrm{lb}}=\int_{M} k_{c}\left(2 H+c_{0}\right)^{2}+\bar{k} K d S \tag{1}
\end{equation*}
$$

where $k_{c}$ and $\bar{k}$ represent specific rigidity constants, $H$ and $K$ are the mean and Gaussian curvatures of the surface $M$, respectively, while $c_{0}$ is the socalled spontaneous curvature.

We would like to mention that this type of generalized bending energy (at times referred to as Helfrich-type energy) has been considered for other types of elastic membranes in biophysics, as well, e.g. [6], [7], [11], [12]. Following the model proposed by Helfrich, two other scientists, S. Choe and X.S. Sun, proposed a similar elastic model for anti-parallel beta sheets, which was published in 2007 in the Biophysical Journal ${ }^{5}$-based on a bending energy, namely:

$$
\begin{equation*}
E_{l \mathrm{~b}}=\int_{M}\left[k\left(H+c_{0}\right)^{2}+\bar{k} K\right] d S, \tag{2}
\end{equation*}
$$

where $d S$ is the infinitesimal area, $c_{0}$ is the "preferred curvature of the surface" as they call it, and $k$ and $\tilde{k}$ are bending moduli that "relate the energy change with changes in mean and Gaussian curvatures".

Our elastic surface model for beta barrels is similar to this model, with the exception of an added constant term that comes from the superficial tension combined with a stress tensor, and with the additional assumption that $c_{0}$ is negligible. We therefore write

$$
\begin{equation*}
E_{b}=\int_{M}\left[k H^{2}+\tilde{k} K+\mu\right] d S, \tag{3}
\end{equation*}
$$



FIGURE 1 Beta Barrel of the GFP (from jelly fish) by scientists Martin Chalfie, Osamu Shimomura and Roger Tsien-Nobel prize recipients (courtesy of the American Association of Clinical Chemistry) [Color figure can be viewed at wileyonlinelibrary.com]


FIGURE 2 Solutions to our boundary value problem, as $H(x)$, and corresponding profile curves $u(x)$ for different $\alpha$ values and for fixed $\epsilon$ value [Color figure can be viewed at wileyonlinelibrary.com]
and call this type of energy generalized Willmore energy (GW energy), or generalized bending energy. The Euler-Lagrange equation corresponding to the GW energy functional can be written as the following (GWE):

$$
\begin{equation*}
\Delta_{g}(H)+2 H\left(H^{2}-K-\epsilon\right)=0 \tag{4}
\end{equation*}
$$

where $\epsilon=\mu / k$ and $\Delta_{g}$ represents the Laplace-Beltrami operator corresponding to the metric $g$ that is naturally induced by the surface parameterization. We are interested in rotational surfaces that represent minimizers of GW energy (that is, solutions to its corresponding EulerLagrange equation described above). These represent our general models for beta barrels. Hereby, we are recalling the basics of a computational study that we have performed on this type of surfaces, in. ${ }^{7}$ Consider a Cartesian system of axes of coordinates $x, y, z$ in $\mathbb{R}^{3}$ and the circles $C_{1}, C_{2}$ of the same radius $\alpha$, centered at $(-1,0,0)$ and ( $1,0,0$ ), situated in planes orthogonal to the $x$ axis. Consider all regular surfaces of revolution of annular-type with boundary $C_{1} \cup C_{2}$. Assume that
among all these surfaces, there exists at least a surface $M$ minimizing the GWE. This surface in assumed embedded in $\mathbb{R}^{3}$ and admitting the representation

$$
M:=\{x, u(x) \cos \varphi, u(x) \sin \varphi\}: x \in[-1,1], \varphi \in[0,2 \pi]
$$

where $u \in C^{4}([-1,1],(0, \infty))$ represents the profile function. Then, the surface $M$ is a solution of the following boundary value problem:

$$
\begin{gather*}
\Delta H+2 H\left(H^{2}-K-\epsilon\right)=0 \quad \text { on } M, \text { where } \epsilon=\frac{\mu}{k}  \tag{5}\\
H=0 \quad \text { on } \partial M=C_{1} \cup C_{2},  \tag{6}\\
u( \pm 1)=\alpha . \tag{7}
\end{gather*}
$$

In these assumptions, there exists a positive value $\alpha^{*}\left(\alpha^{*} \approx 1.5089\right)$ that is independent from the value of $\epsilon$, such that
a If $0<\alpha<\alpha^{*}$, then GWE admits NO minimal solution, that is, any solution satisfies: $H=0$ on $\partial M$ and $H \neq 0$ on $M \backslash(\partial M)$.


FIGURE 3 Solutions to our boundary value problem, as $H(x)$, and corresponding profile curves $u(x)$ for different $\alpha$ values and for fixed $\epsilon$ value [Color figure can be viewed at wileyonlinelibrary.com]
b If $\alpha=\alpha^{*}$, then GWE admits exactly 1 minimal solution (a unique catenoid that exclusively depends on $\alpha^{*}$ ).
c If $\alpha>\alpha^{*}$, then GWE admits exactly 2 minimal solutions ( 2 catenoids whose equations exclusively depend on $\alpha$ ).

The proof is straight-forward, as based on elementary arguments, and can be found in. ${ }^{7}$ Such a Dirichlet boundary value problem was posed for the Willmore equation in [9].

Further, we have analyzed the $\alpha$-family of solutions that corresponds to various fixed values of $\epsilon$. We were able to construct corresponding families of solutions using COMSOL Multiplysics, see fig. 2 and 3. On the other hand, of course, each solution to our boundary value problem (and in particular each catenoidal solution) can actually be represented in COMSOL, if we choose the unique value $\alpha$ appropriately, and deal with the solution branching (in order to graph all corresponding solutions $u$ if that is the case).

Remark the shapes obtained for the profile $u(x)$ as solutions to the boundary value problem associated to GWE, in all the figures presented in this article: they resemble either a catenary, or an undularythus generating catenoidal and unduloidal GW surfaces of revolution. Due to the physical nature of our boundary value problem, the nodoidal solutions are absent, but nodoidal solutions would certainly be present for other types of boundary value conditions of the GWE. Following our analysis, for each and every value of $\alpha$ that is higher than $\alpha^{*}$, there exist 3 distinct solutions, namely 2 catenoidal profiles and a non-minimal solution-which could be unstable (that is, not a local minimizer of the energy). Remark that catenoids represent global minimizers, as Deckelnick and Grunau ${ }^{8}$ showed in a recent article. For the classical Willmore case $\epsilon=0$, authors proved that the non-minimal solution is contained between the 2 catenoids, and it is unstable. Our numerical analysis on the stability of the solutions to the GWE is in progress (Figures 2 and 3 ).

Beta barrels are hereby represented as solutions of the boundary value problem of the GW equation, GWE. Let $D$ represent the "exterior diameter" of the beta barrel, which is defined as the diameter of the end-circles (what we call the imaginary top and bottom of the beta barrel). Let $h$ represent the distance between these two circles, which is informally referred as the height of the beta barrel. Considering the mathematical model of the beta barrel (solution of our boundary value problems), its corresponding height will be 2 (distance between -1 and 1 on the $x$ axis). Correspondingly, the diameter of the "bottom and top circles" will be $2 \alpha$. Therefore, the meaning of the $\alpha$ parameter from the mathematical model is that of the ratio between diameter and height, namely $D / h$, of the real-life beta barrel.

Elasticity theory provides enough reasons for the validity of our computational PDE model. On the other hand, we decided to strengthen our arguments by making an unbiased comparison between our computational model (solution of the boundary value problem) and the statistical models that correspond to protein databases.

## 2 | MATERIALS AND METHODS

## 2.1 | Beta barrel structure dataset from GFP class

To test the performance of the elastic surface model, we used the GFP $C_{\alpha}$ coordinates dataset from Protein Data Bank. Out of the 54 proteins, we excluded 10 with more than 1 domain, and extracted the structure information using Dictionary of Protein Secondary Structure files. We considered only "E" segments that form intradomain $\beta$-ladders as $\beta$-strands.

## 2.2 | Least square estimation using 3 models

We used least square estimation (LSE) to compare the elastic surface model with 2 other models: catenoid and cylinder. In Cartesian system, the coordinates of a catenoid $(x, y, z)$ need to satisfy the following equations:

$$
\left\{\begin{array}{l}
r=a \cosh \left(\frac{z}{a}\right) \\
x=r \cos (2 \pi f) \\
y=r \sin (2 \pi f)
\end{array}\right\}
$$

We used 211 equally spaced $a$ values from 0.4 to 2.5 . For each $a$ value, we generated 300 equally spaced points from -1 to +1 for the variable $z$, and 100 points from 0 to 1 for $f$. Therefore, we obtained 30,000 points for each catenoid with a different $a$ value. The diameter/ height ratio a cosh $\left(\frac{1}{a}\right)$ ranges from 1.5 to 2.7. In Cartesian system, the coordinates of a cylinder ( $x, y, z$ ) need to satisfy the following equations:

$$
\left\{\begin{array}{l}
r \text { is a constant for given } z \\
x=r \cos (2 f \pi) \\
y=r \sin (2 f \pi)
\end{array}\right\}
$$

For the variable $r$, we used 100 equally spaced points, namely from from 0.7 to 1.7. For each $r$ value, we generated 300 equally
spaced points from -1 to +1 for $z$ and, respectively, 100 points from 0 to 1 for $f$. Therefore, we get 30, 000 points for each cylinder with a different $r$ value. The diameter/height ratio $r$ ranges from 0.7 to 1.7. For the elastic surface model, we used 89 equally spaced $\epsilon$ values from 0 to 0.88 , and 48 equally spaced diameter/height values $\alpha$ from 0.7 to 1.17. For each of $\epsilon$ and $\alpha$ combinations, we generated 300 equally spaced points from -1 to +1 for $z$, and 100 points from 0 to 1 for f, respectively. Therefore, by using COMSOL Multiplysics, we obtained 30,000 points for each elastic surface, with different $\epsilon$ and $\alpha$ values.

The procedure to get the LSE for the 3 models can be described as follows:

1. $S$ denotes one simulated dataset. $D$ denotes the $C_{\alpha}$ coordinates data for a certain protein. Denote $n$ as the number of points in dataset $D$.
2. Randomly select $n$ points from $S$, denoted as $s$. Denote $d=D$.
3. Run Procrustes analysis for optimal modeling, up to translating, rotating, and scaling $d$ to superimpose to $s . d$ is updated as the transformed data. The scale value is recorded.
4. Search in dataset $S$ to find $n$ points closest to each of the points in $d$, denoted as $s$.
5. Repeat steps steps 3 ) and 4) 10,000 times, and record $d$ with the least sum of squares for error (SSE) from the subset $s$ in the iteration, denoted as $d_{0}$ and $s_{0}$ with series of scale values as $c_{1}, c_{2}, \ldots, c_{t}$.
6. Denote $d_{1}=d_{0} /\left(c_{1} c_{2} \cdots c_{t}\right)$ as the transformed real data in the original scale. Apply the same translating, rotating, and scaling procedure on $S$ as the one to optimally superimpose $s_{0}$ to $d_{1}$. The result is denoted as $S_{1}$.
7. Apply the Iterative Closest Point method to match $S_{1}$ to $d_{1}$, resulting in rotated data $S_{2}$. Then get correspondences in $S_{2}$ using nearest neighbor search, denoted as $s_{2}$.
8. Calculate $\operatorname{SSE}$ between $s_{2}$ and $d_{1}$.

For each model, we run the above procedure to find the best parameter value to achieve the smallest SSE. The results are shown in the following table.

## 3 | RESULTS AND DISCUSSION

## 3.1 | Statistical model comparison with the computational models in COMSOL

We can see from Table 2 that almost all of the roots of mean squares for error (RMSE) are smaller than 3 , by fitting the 3 models, except proteins 1 fO 9 and 1qyf. It turns out that the 3D structures of these 2 proteins are quite different from the others. Further investigation is needed in order to fit the two proteins using other more proper models. For all the other proteins, the elastic surface model achieves its smallest RMSE when the estimated diameter/height ratio ranges between 0.83 and 0.89 . The cylindrical beta barrel model corresponds

TABLE 2 LSE results for GFP by using 3 models

|  | Catenoid |  |  | Cylinder |  | Elastic surface |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | a | diam./ht. | RMSE | diam./ht. | RMSE | epsilon | diam./ht. | RMSE |
| 1bfp | 0.76 | 1.518 | 2.582 | 0.73 | 1.371 | 0.67 | 0.85 | 1.318 |
| 1c4f | 0.8 | 1.511 | 2.569 | 0.74 | 1.354 | 0.68 | 0.86 | 1.306 |
| 1cv7 | 0.81 | 1.510 | 2.496 | 0.77 | 1.373 | 0.69 | 0.88 | 1.323 |
| 1ema | 0.77 | 1.516 | 2.539 | 0.71 | 1.356 | 0.64 | 0.85 | 1.296 |
| 1 emb | 0.76 | 1.518 | 2.578 | 0.73 | 1.375 | 0.67 | 0.86 | 1.318 |
| 1 emf | 0.79 | 1.512 | 2.527 | 0.72 | 1.359 | 0.63 | 0.85 | 1.293 |
| 1emg | 0.79 | 1.512 | 2.549 | 0.72 | 1.351 | 0.66 | 0.85 | 1.300 |
| 1emk | 0.76 | 1.518 | 2.549 | 0.7 | 1.339 | 0.64 | 0.84 | 1.295 |
| 1 eml | 0.8 | 1.511 | 2.560 | 0.72 | 1.373 | 0.62 | 0.85 | 1.327 |
| 1 emm | 0.76 | 1.518 | 2.540 | 0.72 | 1.380 | 0.66 | 0.87 | 1.322 |
| 1f0b | 0.76 | 1.518 | 2.575 | 0.72 | 1.369 | 0.71 | 0.88 | 1.316 |
| $1 \mathrm{f09}$ | 0.4 | 2.453 | 5.128 | 0.56 | 4.564 | 0.02 | 0.7 | 4.566 |
| 1huy | 0.75 | 1.521 | 2.556 | 0.74 | 1.411 | 0.63 | 0.88 | 1.345 |
| 1jby | 0.77 | 1.516 | 2.590 | 0.74 | 1.361 | 0.67 | 0.87 | 1.321 |
| 1jbz | 0.78 | 1.514 | 2.506 | 0.74 | 1.362 | 0.63 | 0.89 | 1.277 |
| 1 kyp | 0.77 | 1.516 | 2.566 | 0.71 | 1.404 | 0.66 | 0.87 | 1.348 |
| 1 kyr | 0.76 | 1.518 | 2.558 | 0.71 | 1.394 | 0.65 | 0.86 | 1.334 |
| 1 kys | 0.8 | 1.511 | 2.560 | 0.76 | 1.376 | 0.7 | 0.85 | 1.323 |
| 1 mem | 0.75 | 1.521 | 2.532 | 0.75 | 1.361 | 0.65 | 0.86 | 1.292 |
| 1 myw | 0.69 | 1.551 | 2.582 | 0.72 | 1.402 | 0.64 | 0.86 | 1.349 |
| 10xd | 0.74 | 1.525 | 2.598 | 0.72 | 1.387 | 0.68 | 0.86 | 1.337 |
| 10xe | 0.74 | 1.525 | 2.565 | 0.71 | 1.377 | 0.7 | 0.87 | 1.326 |
| 10xf | 0.75 | 1.521 | 2.587 | 0.74 | 1.401 | 0.7 | 0.86 | 1.350 |
| 1q4a | 0.78 | 1.514 | 2.546 | 0.74 | 1.370 | 0.66 | 0.87 | 1.316 |
| 1q4b | 0.75 | 1.521 | 2.571 | 0.72 | 1.372 | 0.64 | 0.88 | 1.322 |
| 1q4c | 0.75 | 1.521 | 2.553 | 0.73 | 1.377 | 0.66 | 0.88 | 1.322 |
| 1q4d | 0.78 | 1.514 | 2.568 | 0.72 | 1.371 | 0.63 | 0.85 | 1.321 |
| 1q4e | 0.76 | 1.518 | 2.549 | 0.74 | 1.380 | 0.66 | 0.86 | 1.332 |
| 1q73 | 0.72 | 1.534 | 2.554 | 0.73 | 1.371 | 0.65 | 0.86 | 1.324 |
| 1qxt | 0.78 | 1.514 | 2.524 | 0.74 | 1.344 | 0.67 | 0.86 | 1.281 |
| 1qy3 | 0.73 | 1.529 | 2.569 | 0.74 | 1.369 | 0.68 | 0.86 | 1.309 |
| 1qyf | 0.4 | 2.453 | 5.108 | 0.51 | 4.022 | 0 | 0.7 | 4.124 |
| 1qyo | 0.79 | 1.512 | 2.554 | 0.76 | 1.419 | 0.68 | 0.86 | 1.350 |
| 1 qyq | 0.78 | 1.514 | 2.526 | 0.7 | 1.412 | 0.66 | 0.87 | 1.342 |
| 1rm9 | 0.78 | 1.514 | 2.513 | 0.74 | 1.406 | 0.63 | 0.88 | 1.323 |
| 1 rmm | 0.77 | 1.516 | 2.536 | 0.72 | 1.387 | 0.67 | 0.87 | 1.318 |
| 1rmo | 0.74 | 1.525 | 2.543 | 0.73 | 1.388 | 0.65 | 0.86 | 1.326 |
| 1 mmp | 0.74 | 1.525 | 2.538 | 0.74 | 1.372 | 0.63 | 0.88 | 1.298 |
|  |  |  |  |  |  |  |  | ntinues) |

TABLE 2 (Continued)

|  | Catenoid |  |  | Cylinder |  | Elastic surface |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | a | diam./ht. | RMSE | diam./ht. | RMSE | epsilon | diam./ht. | RMSE |
| 1rrx | 0.76 | 1.518 | 2.573 | 0.72 | 1.385 | 0.66 | 0.86 | 1.334 |
| 1s6z | 0.79 | 1.512 | 2.598 | 0.71 | 1.418 | 0.67 | 0.86 | 1.365 |
| 2 mmd | 0.76 | 1.518 | 2.537 | 0.74 | 1.402 | 0.6 | 0.87 | 1.321 |
| 2 emn | 0.71 | 1.539 | 2.695 | 0.72 | 1.435 | 0.59 | 0.86 | 1.356 |
| 2 emo | 0.78 | 1.514 | 2.498 | 0.71 | 1.335 | 0.65 | 0.84 | 1.290 |
| 2 yfp | 0.74 | 1.525 | 2.544 | 0.75 | 1.380 | 0.67 | 0.83 | 1.320 |

to a slightly higher RMSE, and similar diameter/height ratios. The catenoidal beta barrel model corresponds to a much higher RMSE and very different estimates of diameter/height ratios.

The "practical reason" for the much better fit of the elastic unduloidal-type model (when compared with its catenoidal counterpart) is the following: On one hand the GFP structure has a small diameter/height ratio, and on the other hand, while both cylinder model and elastic model can fit structures with various diameter/height ratios, the catenoidal model can only allow that ratio to be greater than (or equal to) the value $\alpha^{*}=1.5089$.

The cylindrical model is close enough to the unduloidal-type elastic surface, and definitely more appropriate/accurate than the catenoid. We should not forget that, for the case when $c_{0}$ is not negligible, the GWE equation has $H=c_{0}$ as a stable solution. Therefore, the constant mean curvature (CMC non-zero) model may include circular cylinders, as well as true CMC unduloids. Our elastic model, together with the


FIGURE 4 LSE for GFP 1myw. Black squares represent real data points. Red, green, and blue dots represent the corresponding fitted elastic surface model, catenoid model, and cylinder model

TABLE 3 LSE results for 1cka by using 3 models

|  | Catenoid |  |  | Cylinder |  | Elastic surface |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | a | diam./ <br> ht. | RMSE | diam./ <br> ht. | RMSE | epsilon | diam./ <br> ht. | RMSE |
| 1cka | 1.15 | 1.613 | 0.877 | 1.22 | 1.096 | 0.65 | 1.16 | 1.045 |

statistical analysis, now make it completely clear why the GFP protein model obtained in 2008 by Chalfie, Shimomura and Tsien possesses neither a catenoidal shape, nor a round cylindrical one, but an unduloidal one (see Figure 4). The shape of the profile function $u$ of the beta barrel is always determined by the ratio between the end-circle diameter and height (which is $\alpha$ ). To complete the discussion, we also studied and modeled a protein example 1cka with larger diameter/height ratio. The results in Table 3 show that because the diameter/height ratio is larger for protein 1cka, catenoid model can fit better than cylinder and elastic surface model.

## 4 | CONCLUSION

Our elastic surface model for beta barrels proves itself to be more appropriate than the individual models that were historically used. Physical arguments led us to beta barrel models as solutions to a boundary value problem associated to a general Willmore equation. In addition, our elastic surface model satisfies the requirements of the statistical analysis, as a much better fit than previous models. This is a fortunate case, in which numerical solutions obtained via COMSOL Multiphysics (based on finite element methods) have been tested via statistical analysis, and the results were highly in favor of the elastic model.

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